



Application Note

Description of the phase modulator:

Diffusing an optical channel waveguide at the surface of an X-cut or, more preferably, a Z-cut lithium niobate crystal obtains an optical phase modulator. The main direction of propagation can be oriented parallel to the X- or Y- axis. Cr-Au traveling wave electrodes are deposited over a thick dielectric buffer layer in order to prevent from undesirable optical absorption of the TM-mode, and also in order to get the microwave to optical phase matching condition.

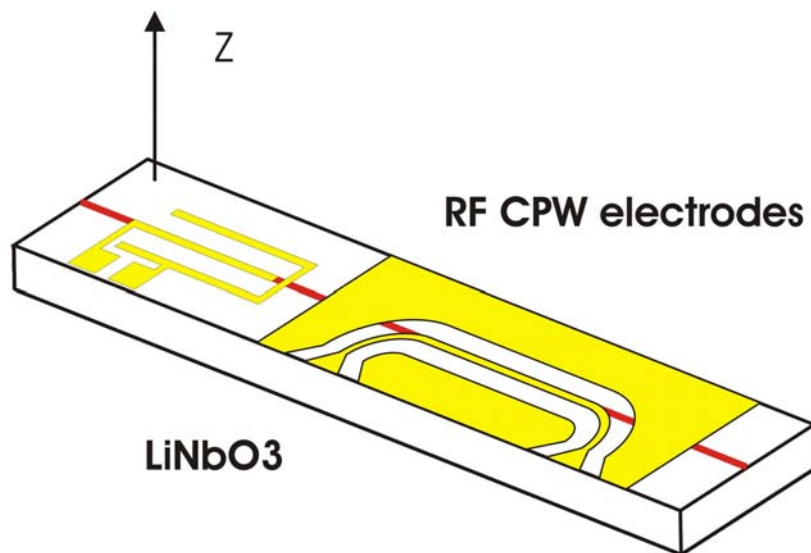


Figure 1: Scheme of a phase modulator on a Z-cut LiNbO3 crystal

Principle of the modulation:

A sinusoidal electrical signal is applied to the electrodes. The amplitude is V_z and the RF frequency pulsation is Ω .

It can be shown that the total phase variation obtained at the output of the modulator illuminated by an optical beam of wavelength λ polarized along the Z-axis, i.e. the TM polarization, is given by the following equation:

$$\phi(t) = \kappa V_z M(\Omega) \cos(\Omega t - \psi) \quad [1]$$

with:

$$\kappa = \frac{\pi}{\lambda G} n_e^3 r_{33} \eta L \quad [2]$$

The frequency dependent amplitude of the signal, limited by phase mismatch and microwave losses is defined by $M(\Omega)$:

$$M(\Omega) = e^{-\left(\frac{\alpha L}{2}\right)} \left[\frac{\sinh^2\left(\frac{\alpha L}{2}\right) + \sin^2\left(\frac{\xi L}{2}\right)}{\left(\frac{\alpha L}{2}\right)^2 + \left(\frac{\xi L}{2}\right)^2} \right]^{\frac{1}{2}} \quad [3]$$

$$\text{where } \xi = (n_\mu - n_o) \frac{\Omega}{c}, \quad \alpha = \alpha_o \sqrt{\Omega} \quad [4]$$

Here, the effective microwave index is n_μ , the characteristic impedance is Z_c and the microwave loss is α . The gap of the electrode is G . the length of the electrodes is L . The electro-optic coefficient of the material is r_{33} and the optical refractive index is n_e . We also consider η as the overlap coefficient between the electric field and the optical field.

One can define at any frequency the effective half-wave voltage of a phase modulator to be:

$$V_\pi(\Omega) = \frac{\lambda G}{n_e^3 r_{33} \eta L M(\Omega)} \quad [5]$$

At DC, the expression becomes:

$$V_\pi = \frac{\lambda G}{n_e^3 r_{33} \eta L} \quad [6]$$

Measurement of the half-wave voltage by spectral analysis:

The measurement of the half-wave voltage at any desired working frequency can be carried by spectral analysis, through a scanning Fabry-Perot interferometer acting as a high-resolution optical spectrum analyzer (OSA). We use for this experiment a laser with a narrow line width (e.g. a 1550nm DFB laser)

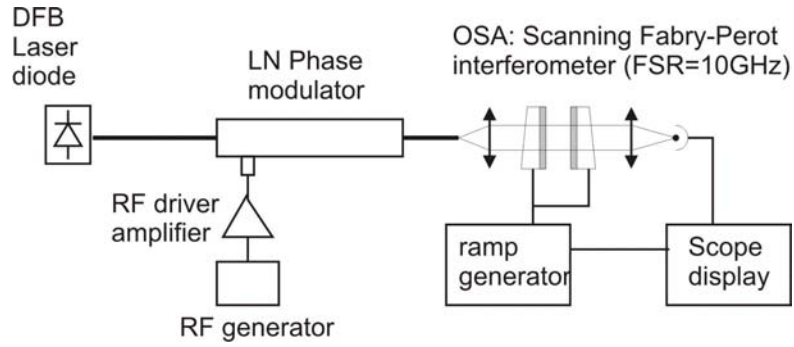


Figure 2: Experimental set-up

We first consider the general expression of the optical field, phase modulated, at the output of the modulator, for a fixed RF input frequency and power:

$$\mathbf{E}(t) = E_0 e^{j\omega t} e^{-j\gamma \cos \Omega t} \quad [7]$$

E_0 is the amplitude of the optical field. ω is the optical pulsation. γ is defined as the index of modulation:

$$\gamma = \kappa \sqrt{Z} M(\Omega) \quad [8]$$

We can now make a spectral decomposition of the signal:

$$\frac{\mathbf{E}(t)}{E_0} = e^{j\omega t} \left[J_0(\gamma) + 2 \sum_{p=1}^{\infty} (-1)^p J_{2p}(\gamma) \cos(2p\Omega t) + 2j \sum_{p'=1}^{\infty} (-1)^{p'} J_{2p'-1}(\gamma) \cos((2p'-1)\Omega t) \right]$$

In this expression, $J_n(\gamma)$ are the Bessel functions. The following figure gives the evolution of the 2 first orders of this family of function.

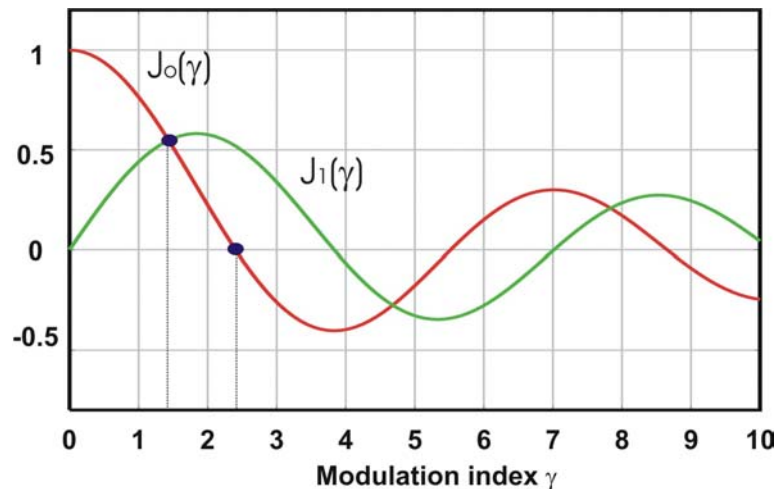


Figure 3: $J_0(\gamma)$ and $J_1(\gamma)$

Hence, the spectral analysis of the optical signal transmitted by a phase modulator driven at a fixed pulsation Ω and with a fixed RF input power yields a spectrum with multiple lines centered around the optical pulsation ω , each line being spaced of Ω from its neighbors. From the signal photo-detected at the output of the interferometer, the amplitude of the central line is given, in relative intensity, by $J_0(\gamma)$, and the two lateral ones by $J_1(\gamma)$.

There are two remarkable points, which correspond to the following situations:

$$J_0(2.4) = 0$$

and

$$J_0(\gamma) = J_1(\gamma) \text{ for } \gamma=1.435$$

These two situations are represented in figures 3 and 4.

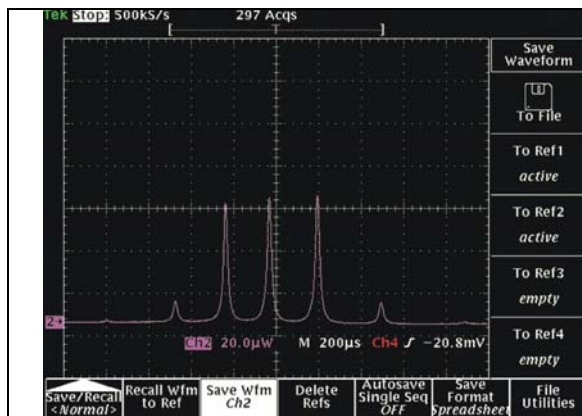


Figure 3: $F=1\text{GHz}$, RF input power:20.5 dBm ($J_0=J_1$)

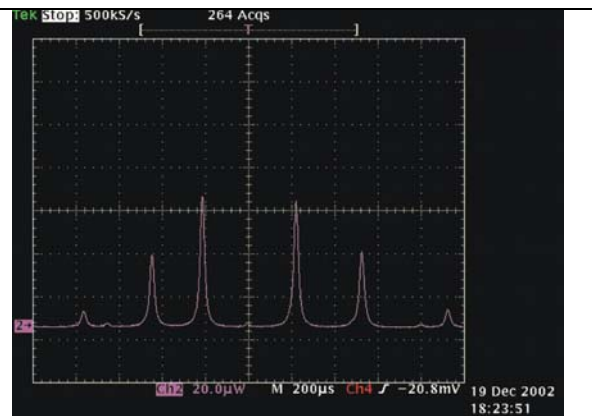


Figure 4: $F=1\text{GHz}$, RF input power:24.6 dBm ($J_0=0$)

We can use this property to calculate accurately the half-wave voltage at the measured frequency:

In the example of the figure, $P=24.6\text{dBm}$ and $f=1\text{GHz}$.

From the RF input power, we get the amplitude of the driving voltage:

$$V_Z=5.37\text{V}$$

We can then deduce the half wave voltage by replacing the modulation index by its value $\gamma=2.4$:

$$V_{\pi}(1\text{GHz}) = \frac{\pi}{2.4} V_Z = 7\text{V}$$